

Consistency of the Douglas – Rachford splitting algorithm for the sum of three nonlinear operators: application to the Stefan problem in permafrost soils

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Abstract: Consistency of the Douglas – Rachford dimensional splitting scheme is proved for the sum of three nonlinear operators constituting an evolution equation. It is shown that the operators must be densely defined, maximal monotone and single valued on a real Hilbert space in order to satisfy conditions, under which the splitting algorithm can be applied. Numerical experiment conducted for a three-dimensional Stefan problem in permafrost soils suggests that the Douglas – Rachford scheme produces reasonable results, although the convergence rate remains unestablished.

Keywords: Splitting Algorithms, Alternating Directions Scheme, Stefan Problem, Nonlinear Heat Equation, Consistent Approximation

1. Introduction

Formulation of the splitting and alternating directions implicit methods (ADI methods) in late 1950-s – early 1960-s [1 - 9] was followed by its rapid developments [10 - 14] and applications in various fields ranging from partial differential equations (PDE) [15, 16] to complex optimization problems [17 – 19]. The reason for such a wide application of these methods is related with a possibility to capture different aspects of a studied system, which are reflected in a complicated PDE, solved by the splitting methods. The complexity of a particular PDE may be due to the presence of terms (operators) that are mathematically very different, making this PDE hard to analyze. In this case, splitting methods provide one with a possibility to split the equation into a set of sub-equations, where each sub-equation is of a type, for which simple and efficient methods are available. The main two ideas for splitting of PDEs are the following (see, e.g. [20]):

1. Physical splitting (related to the underlying processes).
2. Mathematical splitting (related to abstract matrices, found in the differential equations after spatial discretization).

The overall numerical method is then formed by

choosing an appropriate numerical scheme for each sub-equation of the initial PDE and combining the schemes together by operator splitting [21].

Following the classification proposed by Marchuk [11], the factorization and ADI methods, based on inhomogeneous approximation of auxiliary (intermediate) steps, will be called the splitting methods. Under an inhomogeneous finite difference scheme we imply a scheme, for which the coordinate shifts do not induce grid functions from the space, on which the difference operators of the scheme are defined (more details and formal notion of such schemes are developed in [22]). Thus, each intermediate step not necessarily approximates the original problem, but in the whole the approximation takes place. Indeed, the Douglas - Rachford splitting algorithm [6, 23], the study of which is presented in this paper, comprises three steps, each of which separately does not approximate the original problem:

$$(1 + \lambda A_1)(1 + \lambda A_2)(1 + \lambda A_3)U^{n+1} = (\lambda(A_1 + A_2 + A_3) + (1 + \lambda A_1)(1 + \lambda A_2)(1 + \lambda A_3))U^n \quad (1)$$

or, the same, 1 step:

$$U^* = (1 + \lambda A_1)^{-1}(\lambda(A_1 + A_2 + A_3) + (1 + \lambda A_1)(1 + \lambda A_2)(1 + \lambda A_3))U^n \quad (2)$$

2 step:

$$U^{**} = (1 + \lambda A_2)^{-1}(1 + \lambda A_1)^{-1}(\lambda(A_1 + A_2 + A_3) + (1 + \lambda A_1)(1 + \lambda A_2)(1 + \lambda A_3))U^n \quad (3)$$

3 step:

$$U^{n+1} = (1 + \lambda A_3)^{-1}(1 + \lambda A_2)^{-1}(1 + \lambda A_1)^{-1}(\lambda(A_1 + A_2 + A_3) + (1 + \lambda A_1)(1 + \lambda A_2)(1 + \lambda A_3))U^n \quad (4)$$

where λ is the iteration parameter, $A_{i=\{1,2,3\}}$ are some operators and n is the iteration index. From (2) - (4) one can see that the intermediate quantities U^* and U^{**} do not belong to the space of functions that approximate solution of the original problem, while U^n and U^{n+1} are the consecutive iterations of U . To conclude this discussion of relation between ADI and splitting methods we note that in [24, 19] it was shown that the Douglas – Rachford splitting method for minimization of the sum of two monotone operators is a special case of proximal point algorithm.

It is known that the algorithm (1) is absolutely stable if the operators $A_{i=\{1,2,3\}}$ are linear [6, 23]. There has been made numerous attempts to generalize the Douglas - Rachford, Peaceman - Rachford and other splitting algorithms to the case of nonlinear problems [25 – 31]: in [31] a solution-set characterization is used for the estimation of convergence rate of Douglas – Rachford algorithm, applied to variational inequalities and minimization problem with the sum of two convex functions; in [27] the operator-theoretical approach is employed for the studies of convergence of Lie and Peaceman – Rachford splitting, applied to quasilinear parabolic problems with the sum of two nonlinear operators; in [32] the weak convergence of the Douglas – Rachford algorithm was proved for the minimization problem with the sum of two general maximal monotone operators in infinite dimensional spaces; in [33] weak convergence has been proved for an abstract family of projective splitting algorithms for sums of arbitrary numbers of maximal monotone operators (Propositions 3.2 and 4.2 in [33]), but no convergence rate has been established. It has also been demonstrated that the problem with the analysis of dimensional splitting algorithms for the solution of nonlinear differential equations (or equations with variable coefficients) is due to non-commutativity of operators [25]. Despite all these attempts, to the best of our knowledge there is no complete convergence analysis and convergence rate estimation for the problems, comprising the sum of three nonlinear operators available in the literature.

In this paper we present the study of consistency of the algorithm (1) (in application to evolution equations) for the

case when operators $A_{i=\{1,2,3\}}$ are non-linear. For this purpose we employ the results published in [17, 34, 35] and make some additional assumptions that are natural from the point of view of application of the algorithm (1) to Stefan problem in permafrost soils.

The paper is organized as follows: next section contains the theoretical framework and the list of assumptions; in Section 3 the proof of consistency for the algorithm (1) in application to evolution equations is presented; Section 4 is dedicated to the application of algorithm (1) to Stefan problem in permafrost soils. Section 5 presents some numerical results and is followed by conclusions.

2. Notation and Setup

2.1. Theoretical Framework for the Consistency Studies

Let us consider an evolution equation of the form:

$$\begin{aligned} \partial u(t)/\partial t + Au(t) &= 0, \\ u(0) &= u_0, 0 < t < \infty \end{aligned} \quad (5)$$

where H is a Hilbert space and $A: D(A) \subseteq H \rightarrow H$.

From the theory of nonexpansive semigroups it is known that [36, 37 Corollary 31.1]

Proposition 2.1: If the operator A is maximal accretive and $u = u(t)$ is the unique solution of Eq.(5), then

$$u(t) = S(t)u_0 \quad (6)$$

where $u_0 \in D(A)$ and $\{S(t)\}$ is a nonlinear nonexpansive semigroup on $D(A)$, which can be uniquely extended to a nonexpansive semigroup on $D(A)$. The generator of $\{S(t)\}$ on $D(A)$ is $-A$.

Let us remind that the maximal accretive operator is defined as follows:

Definition 2.1: Let $A: D(A) \subseteq H \rightarrow H$ be an operator on the real Hilbert space H . Operator A is called maximal accretive if and only if:

For any $\mu > 0$: $(I + \mu A): D(A) \rightarrow H$ is injective and $(I + \mu A)^{-1}$ is nonexpansive on H .

In our analysis we will make use of the following result [38]:

Proposition 2.2: Let $A: D(A) \subseteq H \rightarrow H$ be an operator on the real Hilbert space H . Then the following two properties of operator A are equivalent:

- (i) A is maximal accretive
- (ii) A is maximal monotone

The maximal monotone operator is defined as follows:

Definition 2.2: Let $A: D(A) \subseteq H \rightarrow H$ be an operator on the real Hilbert space H provided with an inner product (\cdot, \cdot) . Operator A is called maximal monotone if and only if the following two conditions hold:

- (i) $(Au - Av, u - v) \geq 0$ for any $u, v \in D(A)$
- (ii) If $(b - Av, u - v) \geq 0$ for any $v \in D(A)$, then $Au = b$.

At this point we are ready to formulate the main tool of

our further analysis of consistency [34 Corollary 4.3, 35 Theorem 4.3]:

Theorem 2.1: Let $A:D(A) \subseteq H \rightarrow H$ be a maximal accretive densely defined operator and let $\{S(t)\}$ be a semigroup generated by $-A$. Let $\{F(v)\}$ be a family of contractions with the Lipschitz constant $M(v)$ defined on a closed convex subset $C = D(A) \subseteq H$ for all $v > 0$.

If

$$(i) \quad M(v) = 1 + o(v) \text{ for } v \rightarrow 0 \quad (7)$$

$$(ii) \quad \lim_{v \rightarrow 0} (I + \frac{\lambda}{v}(I - F(v)))^{-1}x = (I + \lambda A)^{-1}x \quad (8)$$

for $\forall x \in D(A) \cap C$ and $\forall \lambda > 0$.

Then, for $\forall x \in D(A) \cap C$, the following result holds:

$$\lim_{n \rightarrow +\infty} F(\lambda = t/n)^n u_0 \rightarrow S(t) u_0 \quad (9)$$

And the limit is uniform on bounded intervals of $t \in [0, +\infty)$.

Proof of this theorem can be found in [34, 35].

Another important result that we will employ in our study is the following Lemma proved in [17]:

Lemma 2.1: Let $A:D(A) \subseteq H \rightarrow H$ be a maximal monotone operator. If for every $u \in D(A)$ and $u_\lambda \in H$ there exists a limit

$$\lim_{\lambda \rightarrow 0} (u_\lambda - u)/\lambda = y, \quad (10)$$

Then

$$\lim_{\lambda \rightarrow 0} \frac{u_\lambda - (I + \lambda A)^{-1}u}{\lambda} = P_{Au}(y) \quad (11)$$

where $P_{Au}(y)$ denotes the projection operator for vector y onto the range of A .

The results presented above constitute the basis for the application of nonlinear semigroup theory to the consistency analysis of splitting schemes. In the next Subsection we will make some additional assumptions that are necessary for the specific case of dimensional splitting in the Douglas – Rachford scheme and in Section 3 we will use the Theorem 2.1 and Lemma 2.1 in order to study consistency of this algorithm.

2.2. Finite Differences Formalism for the Quasilinear Heat Equation

The equation describing heat transfer in a system with phase transition has the following form:

$$C(u) \frac{\partial u}{\partial t} = \text{div}(k(u) \text{grad}(u)) \quad (12)$$

where $u = u(x, y, z, t)$ is the unknown function (temperature: $u(t_f) \in C_+^2(\bar{\Omega})$ at every fixed t_f , $C_+^2(\bar{\Omega})$ being a real linear space of positive valued functions), $C(u)$ is the heat capacity, $k(u)$ is thermal conductivity. By Ω we denote a bounded, open subset of an Euclidian space E^3 with boundary $\hat{\Omega}$, the closure of Ω being

denoted by $\bar{\Omega}$.

Following the classification given in [39], we call (12) a quasilinear heat equation. The initial and boundary conditions are taken to be:

$$u(x, y, z, t = 0) = f_0(x, y, z) | (x, y, z) \in \Omega \quad (13)$$

$$u(x, y, z, t) = f(x, y, z, t) | (x, y, z) \in \hat{\Omega}, t \in R_+ \quad (14)$$

We consider the cases when $C(u) \geq \mu > 0$ and $k(u) \geq \kappa > 0$, thus (12) is uniformly parabolic and has a unique solution to the initial-boundary-value problem (13)–(14) [40].

For the formulation of finite difference scheme we introduce the following discretization procedure. Let the vector $h = (h_x, h_y, h_z) \in E^3$ have positive coordinates and G_h be the set of all points $(i \cdot h_x, j \cdot h_y, k \cdot h_z) \in E^3$, the indices i, j, k being integers. Two points $x_1 = (i_1 \cdot h_x, j_1 \cdot h_y, k_1 \cdot h_z)$ and $x_2 = (i_2 \cdot h_x, j_2 \cdot h_y, k_2 \cdot h_z)$ belonging to G_h are called neighbors if $\sqrt{(i_1 - i_2)^2 \cdot h_x^2 + (j_1 - j_2)^2 \cdot h_y^2 + (k_1 - k_2)^2 \cdot h_z^2} = h_{x,y,z}$. The points $x \in G_h \cap \Omega$, all of whose neighbors belong to $\bar{\Omega}$, are denoted by Ω_h .

Following [41], we denote the points $x \in G_h \setminus \Omega_h$ with the property that at least one neighbor belongs to Ω_h by $\hat{\Omega}_h$. Thus, the full spatial mesh is $\bar{\Omega}_h = \Omega_h \cup \hat{\Omega}_h$, $\hat{\Omega}_h$ being the set of boundary points (outside of Ω_h). For the following we assume that $\bar{\Omega}_h$ is a homogeneous (i.e., h_x, h_y, h_z are constants) cubic domain.

Let $T = [0, \infty)$ and $T^0 = T \setminus \{0\}$. Then the time mesh is defined as follows:

$$T_\tau = \{t \in T, t = m \cdot \tau, n = 0, 1, \dots\} \quad (15)$$

The approximate solutions of (12), $U(t)$, are defined on T_τ and take their values in a real finite-dimensional linear space $C(\bar{\Omega}_h)$, the dimension of which is equal to the number of points in $\bar{\Omega}_h$. A function $U(t) \in C(\bar{\Omega}_h)$ is called admissible if $U(t = 0) = f_0(x, y, z)$ and $U = u(x, y, z, t)$ on $\hat{\Omega}_h \times T_\tau^0$. As was pointed out in [41], $u(x, y, z, t)$ is, in general, not known in $\hat{\Omega}_h \times T_\tau^0$ and thus the following assumption should be made: there exists a null sequence (that converges to zero) $\{h_{x,y,z}^\beta\}$ of mesh spacings such that $\Omega_{h^\beta} \subset \Omega$ and $h_{x,y,z}$ always belongs to the sequence $\{h_{x,y,z}^\beta\}$. This assumption guarantees that an admissible function is uniquely specified in $(G_h \times T_\tau^0) \setminus (\Omega_h \times T_\tau^0)$ in terms of the initial and boundary values of $u(x, y, z, t)$.

For every $U(t) \in C(\bar{\Omega}_h)$ we define the first-order forward and backward difference operators:

$$\begin{cases} h_\alpha \cdot D_\alpha^+ U(t) = U_{i_\alpha+1}(t) - U_{i_\alpha}(t) \\ h_\alpha \cdot D_\alpha^- U(t) = U_{i_\alpha-1}(t) - U_{i_\alpha}(t) \end{cases} \quad (16)$$

where $\alpha = \{x, y, z\}$, $2 \leq i_\alpha \leq (\text{dim}_\alpha(\bar{\Omega}_h) - 1)$, $\text{dim}_\alpha(\bar{\Omega}_h)$ being the number of points in $\bar{\Omega}_h$ along α -direction in the Cartesian coordinates. For any suitably defined function $U(t) \in C(\bar{\Omega}_h)$ we set

$$k_{i_\alpha \pm 1/2, n} = k\left(\frac{U_{i_\alpha \pm 1} + U_{i_\alpha}}{2}\right) \tag{17}$$

Thus, $k_{i_\alpha \pm 1/2, n}$ are the values of $k(U(t))$ at the fictitious intermediate nodes of the mesh Ω_h at $t = \tau \cdot n \in T_\tau$, $U(t) \in C(\bar{\Omega}_h)$.

From the Eqs. (16)--(17) one obtains:

$$-D_{i_\alpha}^+(k_n(U) \cdot D_{i_\alpha}^- U) = \frac{1}{h_{i_\alpha}}(\Lambda_{i_\alpha, n}^+ + \Lambda_{i_\alpha, n}^-)U, \tag{18}$$

where

$$\begin{cases} \Lambda_{i_\alpha, n}^- U = k_{i_\alpha - 1/2, n} \cdot \frac{U_{i_\alpha - 1} - U_{i_\alpha}}{h_{i_\alpha}} \\ \Lambda_{i_\alpha, n}^+ U = k_{i_\alpha + 1/2, n} \cdot \frac{U_{i_\alpha + 1} - U_{i_\alpha}}{h_{i_\alpha}} \end{cases} \tag{19}$$

and $k_n(U)$ corresponds the value of $k(U(t))$ at $t = \tau \cdot n \in T_\tau$, $U(t) \in C(\bar{\Omega}_h)$, i_α and α being defined as $\alpha = \{x, y, z\}$; $1 \leq i_\alpha \leq \dim_\alpha(\bar{\Omega}_h)$.

We also provide $C(\bar{\Omega}_h)$ with the l_2 -inner product (U, V) and induced norm $\|U\| = (U, U)^{1/2}$ on $\bar{\Omega}_h$:

$$(U, V) = h_x \cdot h_y \cdot h_z \sum_{p \in G_h} U(p)V(p) \tag{20}$$

The maximum norm is defined as following:

$$\|U\|_\infty = \max(U) \tag{21}$$

We define on $\bar{\Omega}_h$ a modified A -inner product $(U, V)_{n, A}$ and induced norm

$$\|U\|_{n, A} = (U, U)_{n, A}^{1/2}$$

by

$$(U, V)_{n, A} = h_x \cdot h_y \cdot h_z \sum_{\alpha = \{x, y, z\}} \sum_{p \in G_h} U(p)A_{n, \alpha}V(p) \tag{22}$$

$$A_{n, \alpha}(V(p)) = -C_n^{-1}(V(p)) \cdot D_\alpha^+(k_n(V(p)))D_\alpha^-(V(p)) \tag{23}$$

where $C_n^{-1}(V(p))$ and $k_n(V(p))$ correspond the values of $C^{-1}(v(x, y, z, t))$ and $k(v(x, y, z, t))$ at $p = (x, y, z) \in G_h, t = \tau \cdot n \in T_\tau$, $v(t) \in C(\bar{\Omega})$ respectively. Note that the so defined operator $A_{n, \alpha}$ implies that it is generally not true that $A_{n, \alpha}: C(\bar{\Omega}_h) \rightarrow C(\bar{\Omega}_h)$.

Now we can write the following Douglas – Rachford finite differences scheme for the heat equation (see detailed discussion of this scheme in [42]):

$$\begin{aligned} & (1 + \tau \cdot A_{n, x})(1 + \tau \cdot A_{n, y})(1 + \tau \cdot A_{n, z})U_\tau^n \\ & = \sum_\alpha A_{n, \alpha}(U^n) \end{aligned} \tag{24}$$

where $A_{n, \alpha}$ are given by (23), $\alpha = \{x, y, z\}$, $U_\tau^n = (U^{n+1} - U^n)/\tau$, $U^n = U(t)$ and $U^{n+1} = U(t + \tau)$ at $t = n \cdot \tau \in T_\tau$, $U(t) \in C(\bar{\Omega}_h)$.

Application of the scheme (24) to the Stefan problem will be studied in Section 4.

To complete this Section we present the following two

assumptions that are necessary to guarantee computational stability of this scheme [42]

Assumption 2.1: In the scheme (24) the operators $A_{n, \alpha}$ are acting on the vectors U^{n+1} in the following way:

$$A_{n, \alpha}(U^{n+1}) = -C_n^{-1}(U^n) \cdot D_\alpha^+(k_n(U^n))D_\alpha^-(U^{n+1}) \tag{25}$$

This assumption leads to consideration of a linear finite-difference scheme of type (24) for the heat equation with variable coefficients C_n and k_n , the forms of which depend on the temperature field $U^n = U(t)$ calculated at the previous time level. This kind of assumption is commonly used when considering nonlinear time dependent problems (e.g., [26]).

Let us rewrite (25) in the following equivalent form:

$$A_{n, \alpha}(V(p)) = \bar{A}_{n, \alpha}(V(p)) + \delta_{n, \alpha}(V(p)) \tag{26}$$

where

$$\begin{cases} \bar{A}_{n, \alpha}(V(p)) = \frac{-K_{i_\alpha, n}^+ D_\alpha^+(D_\alpha^- V(p))}{C_n} \\ \delta_{n, \alpha}(V(p)) = \frac{K_{i_\alpha, n}^-(D_\alpha^+ - D_\alpha^-)V(p)}{C_n} \end{cases} \tag{27}$$

$$\begin{cases} K_{i_\alpha, n}^+ = (k_{i_\alpha + \frac{1}{2}, n} + k_{i_\alpha - \frac{1}{2}, n})/2 \\ K_{i_\alpha, n}^- = (k_{i_\alpha + \frac{1}{2}, n} - k_{i_\alpha - \frac{1}{2}, n})/(2 \cdot h_\alpha) \end{cases} \tag{28}$$

Assumption 2.2: The functions $C^{-1}: C(\bar{\Omega}_h) \rightarrow \ell^\infty(\bar{\Omega}_h)$, $K^+: C(\bar{\Omega}_h) \rightarrow \ell^\infty(\bar{\Omega}_h)$ and $K^-: C(\bar{\Omega}_h) \rightarrow \ell^\infty(\bar{\Omega}_h)$ (defined by (17) and (28)) are mapping the elements $U^n(p) \in C(\bar{\Omega}_h)$ into a sequence space $\ell^\infty(\bar{\Omega}_h)$, the elements of which $(K_n^+(p), K_n^-(p), C_n^{-1}(p) \in \ell^\infty(\bar{\Omega}_h))$ satisfy the following relations:

$$\begin{cases} C^{-1}(U^n(p)) = C_n^{-1}(p) \\ |D_\alpha^+(C_n^{-1}(p))| \leq c_1 \cdot C_n^{-1}(p) \\ |D_\alpha^-(C_n^{-1}(p))| \leq c_1 \cdot C_n^{-1}(p) \end{cases} \tag{29}$$

$$\begin{cases} K^+(U^n(p)) = K_n^+(p) = \frac{k_{i_\alpha + \frac{1}{2}, n}(p) + k_{i_\alpha - \frac{1}{2}, n}(p)}{2} \\ |D_\alpha^+(K_n^+(p))| \leq c_2 \cdot K_n^+(p) \\ |D_\alpha^-(K_n^+(p))| \leq c_2 \cdot K_n^+(p) \end{cases} \tag{30}$$

$$\begin{cases} K^-(U^n(p)) = K_n^-(p) = \frac{k_{i_\alpha + \frac{1}{2}, n}(p) - k_{i_\alpha - \frac{1}{2}, n}(p)}{2} \\ |D_\alpha^+(K_n^-(p))| \leq c_3 \cdot K_n^-(p) \\ |D_\alpha^-(K_n^-(p))| \leq c_3 \cdot K_n^-(p) \end{cases} \tag{31}$$

where the operators D_α^+ , D_α^- are defined similarly to (16); c_1, c_2, c_3 are real positive constants. $\ell^\infty(\bar{\Omega}_h)$ is the Banach space with the maximum norm.

3. Consistency of the Douglas – Rachford Splitting Scheme

In order to proceed further and consider the case where the operator A in (5) is the sum of three nonlinear operators,

$$A = A_1 + A_2 + A_3 \quad (32)$$

we have to make the following assumptions.

Assumption 3.1: The operators

$$A_{i=\{1,2,3\}}: D(A_{i=\{1,2,3\}}) \subseteq H \rightarrow H \quad (33)$$

and

$$A = A_1 + A_2 + A_3$$

are single valued, densely defined and maximal monotone on a Hilbert space H .

Assumption 3.2: The range of $I - \lambda(A_1 + A_2 + A_3)$ is dense in H for all $\lambda > 0$.

These two assumptions lead to the following important result:

Lemma 3.1: If Assumptions 3.1 and 3.2 are valid, then

$$\overline{\text{graph}(A_1 + A_2 + A_3)} = \text{graph}(A)$$

The proof of this Lemma is presented in [27].

Now we are ready for the formulation of the main Theorem of consistency analysis for the Douglas – Rachford scheme.

Theorem 3.1: Let the Assumption 3.1 and Assumption 3.2 hold. Then, the family of contractions

$$F(\lambda) = J_{A_3}^\lambda J_{A_2}^\lambda J_{A_1}^\lambda (\lambda(A_1 + A_2 + A_3)) + I \quad (34)$$

defined on a closed convex subset $C = D(A) \subseteq H$ for all $\lambda > 0$ with

$$J_{A_{i=\{1,2,3\}}}^\lambda = (I + \lambda A_{i=\{1,2,3\}})^{-1}, \quad (35)$$

satisfies the conditions (i) and (ii) of the Theorem 2.1 and $-A = -(A_1 + A_2 + A_3)$ generates a semigroup $\{S(t)\}$ such that $u(t) = S(t)u_0$ and

$$\lim_{n \rightarrow +\infty} F(\lambda = t/n)^n u_0 \rightarrow S(t)u_0$$

holds for the evolution equation (5).

Proof: Satisfaction of the first condition of the Theorem 2.1 is obvious when taking the limit $\lim_{\lambda \rightarrow 0} F(\lambda)$.

For the proof of satisfaction of the second condition of the Theorem 2.1 let us consider the following quantity:

$$\begin{aligned} \frac{1}{\lambda}(U_\lambda - F(\lambda)U_\lambda) &= \frac{1}{\lambda}(J_{A_3}^\lambda J_{A_2}^\lambda J_{A_1}^\lambda ((J_{A_1}^\lambda)^{-1} + (J_{A_2}^\lambda)^{-1} \\ &\quad + (J_{A_3}^\lambda)^{-1} - 3)U_\lambda) \end{aligned} \quad (36)$$

Note that in general case when the operators $A_{i=\{1,2,3\}}$ are nonlinear, $J_{A_i}^\lambda$, $(J_{A_i}^\lambda)^{-1}$ and $J_{A_j}^\lambda$ do not commute with each other. At the same time, if $A_{i=\{1,2,3\}}$ are single valued, one has the following identity:

$$(I + \lambda A_{i=\{1,2,3\}})J_{A_{i=\{1,2,3\}}}^\lambda = I \quad (37)$$

Thus, we have for (36):

$$\begin{aligned} \frac{1}{\lambda}(U_\lambda - F(\lambda)U_\lambda) &= \frac{1}{\lambda}(J_{A_3}^\lambda J_{A_2}^\lambda + J_{A_3}^\lambda J_{A_1}^\lambda + J_{A_2}^\lambda J_{A_1}^\lambda \\ &\quad - 3J_{A_3}^\lambda J_{A_2}^\lambda J_{A_1}^\lambda + J_{A_3}^\lambda J_{A_2}^\lambda [J_{A_1}^\lambda, (J_{A_2}^\lambda)^{-1}] + \\ &\quad + J_{A_3}^\lambda [J_{A_2}^\lambda, (J_{A_3}^\lambda)^{-1}]J_{A_1}^\lambda + J_{A_3}^\lambda J_{A_2}^\lambda [J_{A_1}^\lambda, (J_{A_3}^\lambda)^{-1}])U_\lambda \end{aligned} \quad (38)$$

where

$$[J_{A_i}^\lambda, (J_{A_j}^\lambda)^{\pm 1}] = J_{A_i}^\lambda (J_{A_j}^\lambda)^{\pm 1} - (J_{A_j}^\lambda)^{\pm 1} J_{A_i}^\lambda \quad (39)$$

Proceeding in the same manner we obtain the following expression:

$$\begin{aligned} \frac{1}{\lambda}(U_\lambda - F(\lambda)U_\lambda) &= \frac{1}{\lambda}(J_{A_2}^\lambda J_{A_1}^\lambda (U_\lambda - J_{A_3}^\lambda U_\lambda) + \\ &\quad J_{A_3}^\lambda J_{A_1}^\lambda (U_\lambda - J_{A_2}^\lambda U_\lambda) + \\ &\quad J_{A_3}^\lambda J_{A_2}^\lambda (U_\lambda - J_{A_1}^\lambda U_\lambda) + \\ &\quad J_{A_3}^\lambda [J_{A_1}^\lambda, J_{A_2}^\lambda]U_\lambda + J_{A_3}^\lambda J_{A_2}^\lambda [J_{A_1}^\lambda, (J_{A_2}^\lambda)^{-1}]U_\lambda + \\ &\quad J_{A_3}^\lambda [J_{A_2}^\lambda, (J_{A_3}^\lambda)^{-1}]J_{A_1}^\lambda U_\lambda + \\ &\quad J_{A_3}^\lambda J_{A_2}^\lambda [J_{A_1}^\lambda, (J_{A_3}^\lambda)^{-1}]U_\lambda + \\ &\quad J_{A_2}^\lambda [J_{A_1}^\lambda, J_{A_3}^\lambda]U_\lambda + [J_{A_2}^\lambda, J_{A_3}^\lambda]J_{A_1}^\lambda U_\lambda. \end{aligned} \quad (40)$$

Taking the limit at $\lambda \rightarrow 0$ and applying Lemma 2.1 under the condition

$$\lim_{n \rightarrow \infty} U_{\lambda=t/n} = u(t)$$

we obtain:

$$\lim_{\lambda \rightarrow 0} \frac{1}{\lambda}(U_\lambda - F(\lambda)U_\lambda) = P_{A_1 u}(0) + P_{A_2 u}(0) + P_{A_3 u}(0) \quad (41)$$

Thus, the operators (34) satisfy the conditions (i) and (ii) of the Theorem 2.1. Assumption 3.1 and Assumption 3.2 together with the Proposition 2.2 lead to the conclusion that operators

$$A_{i=\{1,2,3\}}: D(A_{i=\{1,2,3\}}) \subseteq H \rightarrow H$$

and $A = A_1 + A_2 + A_3$ are maximal accretive. Finally, Proposition 2.1 leads to the conclusion that $-A = -(A_1 + A_2 + A_3)$ generates a semigroup $\{S(t)\}$ such that $u(t) = S(t)u_0$ and

$$\lim_{n \rightarrow +\infty} F(\lambda = t/n)^n u_0 \rightarrow S(t)u_0$$

holds for the evolution equation (5). The Proof is

completed.

Corollary 3.1: The algorithm (1) provides a consistent approximation for the evolution equation (5).

Proof: The statement becomes obvious when the algorithm (1) is put in the form:

$$U^{n+1} = F(\lambda)U^n.$$

4. Application to Stefan Problem in Permafrost Soils

4.1. Stefan problem in permafrost soils

Following [42], we introduce some additional notations and notions that are necessary for the study of Stefan problem in permafrost soils.

Basing on results of [43 - 45], we formulate the model for the Stefan problem without explicitly invoking the front-tracking condition. This approach is justified by the observations that in case of explicit front-tracking models applied to Stefan problem in salted permafrost soils, there appears an overcooled region (frozen fringe zone) [46, 47] (the analysis of this zone and related frost heave and cryogenic suction processes [48, 49] are beyond the scope of the present work). Thus, we take into account the phase transition by introducing effective heat capacity $C_{eff}(u)$, which incorporates the latent heat per unit mass L_0 :

$$\begin{cases} C_{eff}(u) = C_0(u) - (1 - w(u)) \cdot \rho_w \cdot L_0 \frac{dw(u-0)}{du}, \\ C_0(u) = C_{th} + (C_f - C_{th}) \cdot w(u), \end{cases} \quad (42)$$

$$w(u) = \begin{cases} 1 - \frac{1}{1+S \cdot (T_{ph}-u)}, & u \leq T_{ph} \\ 0, & u > T_{ph} \end{cases} \quad (43)$$

where C_{th} and C_f are the values of volumetric heat capacities of the soil in thawed and frozen phases respectively, $w(u)$ is the fraction of frozen water, S is the smoothing parameter, T_{ph} is the phase transition temperature, ρ_w is the density of water. The thermal conductivity is taken to be

$$k(u) = k_{th} + (k_f - k_{th}) \cdot w(u) \quad (44)$$

where k_{th} and k_f are the values of thermal conductivities in thawed and frozen phases respectively.

4.2. Consistent Approximation for the Quasilinear Heat Equation

In Section 3 we have shown that the algorithm (1) can be used for a consistent approximation of evolution equation comprising the sum of three nonlinear operators. Let us now show that the scheme (24) provides a consistent approximation for the quasilinear heat equation (12).

Proposition 4.1: Let the Assumption 2.1 hold and let $C_n^{-1}(V(p))$ and $k_n(V(p))$ correspond the values of $C^{-1}(v(x, y, z, t))$ and $k(v(x, y, z, t))$ at $p = (x, y, z) \in$

G_h , $t = \tau \cdot n \in T_\tau$, $v(t) \in C(\bar{\Omega})$ respectively. If $C^{-1}(v) = C_{eff}^{-1}(v)$ (in accordance with (42) - (43)) and $k(v)$ is defined by (43) - (44), then the scheme (24):

$$(1 + \tau A_{n,x})(1 + \tau A_{n,y})(1 + \tau A_{n,z})U_\tau^n = \sum_\alpha A_{n,\alpha}(U^n)$$

gives a consistent approximation for equation (12):

$$C(u) \frac{\partial u}{\partial t} = div(k(u)grad(u)).$$

Proof: The scheme (24) takes the form of the algorithm (1) if we put $\tau = \lambda$ and $A_{n,i=\{x,y,z\}} = A_{n,i=\{1,2,3\}}$. According to the Theorem 3.1 and Corollary 3.1, algorithm (1) provides a consistent approximation for the evolution equation of type (5) if all the operators

$$A_{i=\{1,2,3\}}: D(A_{i=\{1,2,3\}}) \subseteq H \rightarrow H$$

and

$$A = A_1 + A_2 + A_3$$

are

- (i) densely defined,
- (ii) single valued,
- (iii) maximal monotone on a Hilbert space H .

The first property comes from the definition of linear space $C(\Omega_h)$, given in Subsection 2.2. Indeed, $D(A_{i=\{1,2,3\}}) = C(\Omega_h)$ and $H = C(\Omega_h)$.

The second and third properties come from the definition of $C_n^{-1}(V(p))$, $k_n(V(p))$ and from the fact that all the components of an admissible function $U(t) \in C(\Omega_h)$ are positive (see Subsection 2.2).

Thus, the scheme (24) satisfies the consistency property (9). Proof is completed.

In [42] the sufficient criterion for the computational stability of the algorithm (1) has been obtained. Below we present without proof the formulation of the main theorem for the analysis of computational stability of algorithm (1) (the proof can be found in [42]):

Theorem 4.1: Let the operators $A_{n,\alpha}$ (defined by (25)) satisfy the following condition for any $n \in T_\tau$

$$(1 - c_4 \cdot \epsilon^n) \sum_\alpha A_{n-1,\alpha} \leq \sum_\alpha A_{n,\alpha} \leq (1 + c_4 \cdot \epsilon^n) \sum_\alpha A_{n-1,\alpha}, \quad (45)$$

where $\epsilon^n = \|U^n - U^{n-1}\|_\infty$, $c_4 > 0$ is a constant.

Then the relation:

$$(1 + \tau A_{n,x})(1 + \tau A_{n,y})(1 + \tau A_{n,z}) \geq \frac{\tau}{1+\rho} \sum_\alpha A_{n,\alpha}, \quad (46)$$

(with n from (15), $\rho \geq 1$) is sufficient for the following estimate to hold:

$$(\sum_\alpha A_{n,\alpha} U^{n+1}, U^{n+1}) \leq \rho^2 (1 + c_4 \epsilon^n) (\sum_\alpha A_{n-1,\alpha} U^n, U^n), \quad (47)$$

where $U^n, U^{n+1} \in C(\bar{\Omega}_h)$ are the solutions of Eq.(24).

Application of this theorem to the scheme (24) leads to the following stability criterion [42]:

$$\frac{\tau}{h_a^2} \leq \frac{\rho}{1+\rho} \frac{\min(K_n^+)}{\max(C_n)} \left(\frac{\min(C_n)}{\max(K_n^+)} \right)^2 \cdot c_0, \quad (48)$$

where c_0 is a constant. The criterion (48) is sufficient for relation (47) to hold, which reflects the notion of computational stability.

Thus, for the linearized finite differences scheme (24) we have proved the consistency and the stability criterion is obtained in [42]. With these results we can proceed further and analyze the numerical experiments.

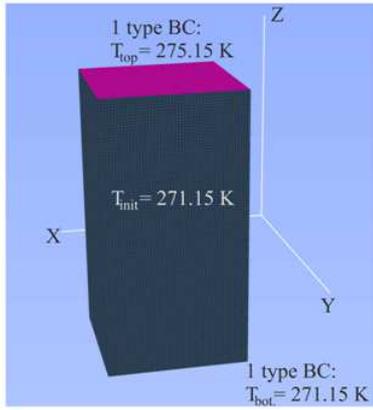


Figure 1. Rectangular parallelepiped with the mesh used in numerical experiment.

5. Numerical Experiment

Let us consider a rectangular parallelepiped with dimensions

$$5[m] \times 5[m] \times 10[m]$$

and following initial and boundary conditions (see Fig. 1):

$$T_{top} = 275.15[K]$$

$$T_{bottom} = 271.15 [K]$$

$$T_{init} = u_0 = 271.15 [K]$$

On the side edges of this parallelepiped zero thermal fluxes are posed.

Let us suppose that the specific heat and thermal conductivity of material in this parallelepiped are given by (42) – (44) with the following values of constants:

$$C_{th} = 1.89 [MJ/(m^3K)],$$

$$C_f = 1.74 [MJ/(m^3K)],$$

$$k_{th} = 0.7 [W/(mK)],$$

$$k_f = 0.75 [W/(mK)],$$

$$T_{ph} = 273.15 [K],$$

$$S = 10[K^{-1}],$$

$$L_0 = 334 [kJ/kg],$$

$$\rho_w = 1000 [kg/m^3]$$

Let us employ a homogeneous mesh with $\Delta x = \Delta y = \Delta z = 0.1 [m]$ and the following number of nodes: $51 \times 51 \times 100$. Using stability criterion (48) we obtain the values $\tau = 1160 [s]$ for the maximal time step and 744 for the minimal number of iterations.

In Fig. 2 the calculation results are presented for the time period of 10 days. Every contour plot in Fig. 2 presents a temperature field in the plane of parallelepiped with the coordinate $x = 7 [m]$.

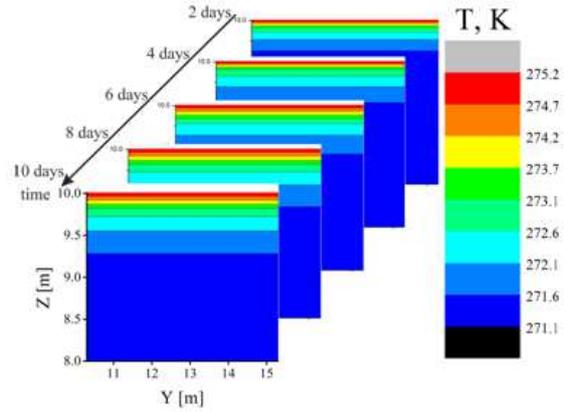


Figure 2. Contour plots of the temperature field at the planes $x = 7 [m]$ obtained with the scheme (24) at different time moments.

One can observe a slight shift of the freezing front into the depth of soil. In order to demonstrate the change better, we plot the dependence of temperature on the coordinate in Fig. 3.

In Fig. 3 the same results as in Fig. 2 are presented in the form of temperature change with the depth along the line with coordinates $x = 7 [m]$, $y = 13.6 [m]$.

One can observe the phase transition point and different behavior of temperature field below and above that point. This is in agreement with Eq. (42) – (44) and the general expectations for the results of solution of a Stefan problem.

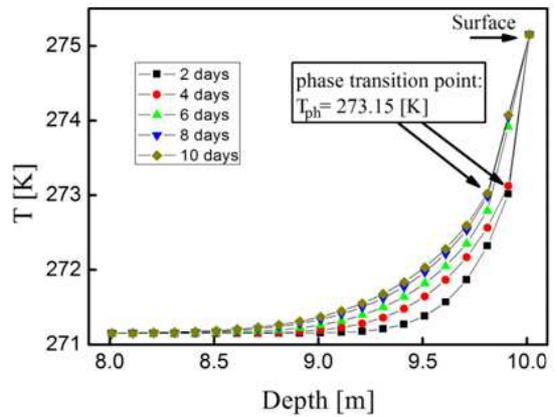


Figure 3. The dependence of temperature on the vertical coordinate (depth) at different time moments.

6. Conclusions

We have shown that the Douglas – Rachford splitting algorithm (1) can be employed for a consistent approximation of evolution equation of type (5) (Theorem 3.1 and Corollary 3.1), provided the operators

$$A_{i=\{1,2,3\}}: D(A_{i=\{1,2,3\}}) \subseteq H \rightarrow H$$

and

$$A = A_1 + A_2 + A_3$$

are single valued, densely defined and maximal monotone in a Hilbert space H .

We suggest that the finite differences scheme (24) gives a consistent approximation for the heat equation (12) and can be used for the study of Stefan problems (Proposition 4.1). The conducted numerical experiment (with the computational stability criterion taken from [42]) provides evidence that the Douglas – Rachford dimensional splitting algorithm can be used for the studies of heat transfer in permafrost soils (we argue that all other experiments (with different parameters) were also successful). Although the numerical experiments conducted with the stability criterion (48) appear to be successful, we note that the convergence rate has not yet been established for the scheme (24) and this presents an open problem.

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