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# Analysis of turbulent Hydromagnetic flow with Radiative heat over a moving vertical plate in a Rotating System

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**Abstract:** In this paper, the combined effects of magnetic fields, buoyancy force, thermal radiation, viscous and Ohmic heating on turbulent hydromagnetic flow of an incompressible electrically conducting fluid over a moving vertical plate in a rotating system is investigated numerically. The governing equations are reduced to non-linear ordinary differential equations using the time-averaged approach known as Reynolds-averaged Navier–Stokes equations (or RANS equations) and tackled by employing an efficient Runge-Kutta Fehlberg integration technique coupled with shooting scheme. Graphical results showing the effects of various thermophysical parameters on the velocity, temperature, local skin friction and local Nusselt number are presented and discussed quantitatively.

**Keywords:** Turbulent Flow, Magnetohydrodynamics (MHD), Radiative Heat, Time Averaging, Rotating System

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## 1. Introduction

Turbulent flow is a flow regime characterized by chaotic variation in the fluid properties, such as low momentum diffusion, high momentum convection, and rapid variation of pressure and velocity in space and time. In turbulent flow, drag due to boundary layer increases and the unsteady vortices appear on many scales and interact with each other [1]. Turbulence causes the formation of eddies of many different length scales. Studies related to turbulent flow and heat transfer not only present a mathematical challenge but find several applications in many industrial, engineering and technological processes [2]. For instance, the external flow over all kind of vehicles such as cars, airplanes, ships and submarine are turbulence. The flow conditions in many industrial equipment such as pipes, ducts, precipitators, gas scrubbers, dynamic scraped surface heat exchangers, internal combustion engines and gas turbines are turbulence. In many geophysical flows such as rivers and atmospheric boundary layer, the flow turbulence is dominated by the coherent structure activities and associated turbulent events. In the medical field of cardiology, a stethoscope is used to detect heart\_sounds and bruits, which are due to turbulent

blood flow. Moreover, when flow is turbulent, particles exhibit additional transverse motion which enhances the rate of energy and momentum exchange between them thus increasing the heat transfer and the friction coefficient [3].

Meanwhile, hydromagnetic flow and heat transfer have received considerable attention in recent years due to its various applications in science, engineering and industries. Melt refining involves magnetic field applications to control excessive heat transfer rate. Other applications of magnetohydrodynamic (MHD) heat transfer include MHD generators, plasma propulsion in astronautics, nuclear reactor thermal dynamics and ionized-geothermal energy systems [4-6]. Model studies of the phenomena of MHD convection with respect to turbulent flow have been made by many authors. Burr *et al.* [7] presented the analysis of hydromagnetic turbulent flow and heat transfer in a rectangular duct with strong sidewall jets. Ji *et al.* [8] conducted a numerical investigation into turbulent pipe flow under the influence of an imposed transverse magnetic field. Kenjere and Hanjalic [9] numerically studied the effects of Lorentz force in turbulence closure models. Kitamura and Hirata [10] investigated the problem of turbulent heat and momentum transfer for electrically

conducting fluid flowing in two-dimensional channel under the influence of an imposed transverse magnetic field. Numerical simulation of large-eddy in conductive flows at low magnetic Reynolds number was reported by Knaepen and Moin [11]. Kobayashi [12] also presented large eddy simulation of hydromagnetic turbulent channel flows with local subgrid-scale model based on coherent structures. Meanwhile, all conducting fluids with a temperature greater than absolute zero emit thermal radiation. Thermal radiation is the emission of electromagnetic waves. It represents a conversion of thermal energy into electromagnetic energy [13]. Thermal energy results in kinetic energy in the random movements of molecules in the conducting fluid. Therefore, whenever the temperature of surrounding fluid is high, the radiation effects play a very important role in the flow process [14]. This situation does occur in many engineering and industrial flow systems. In such cases one has to take into account the effects of radiation and free convection. Aboeldahab and Gendy [15] studied the radiation effects on MHD free convective flow of a gas past a semi-infinite vertical plate with variable thermophysical properties for higher-temperature difference. Ishak [16] investigated the thermal radiation effects on hydro-magnetic flow due to an exponentially stretching sheet. Furthermore, hydromagnetic flow in a rotating environment as received significant attention of several researchers due to its applications in various technological situations which are governed by the action of Coriolis force. Oceanography, meteorology, atmospheric science and limnology all contain some important and essential features of rotating fluids. An order of magnitude analysis shows that in basic field equations the effects of Coriolis force is more significant as compared to that of inertia and viscous forces. It is worthy to note that Coriolis and magnetohydrodynamic forces are comparable in magnitude and Coriolis force induces secondary flow in the fluid. Seth *et al.* [17] investigated unsteady hydromagnetic Couette flow of a viscous incompressible and electrically conducting fluid in a rotating system in the presence of uniform transverse magnetic field considering different aspects of the problem. Kinyanjui *et al.* [18] investigated the effect of hall current on MHD turbulent flow over a vertical plate in a rotating system. From the literature review, it appears that the combined effects of magnetic fields, buoyancy force, thermal radiation, viscous and Ohmic heating on turbulent hydromagnetic boundary layer flow in a rotating system has not been reported.

The main objective of this present study is to extend the theoretical model of MHD turbulent flow in a rotating system as described in [18] to include the combined effects of magnetic fields, buoyancy force, thermal radiation, viscous and Ohmic heating. The mathematical formulation of the problem based on the Reynolds averaged Navier-Stokes (or RANS) method is established and shooting method coupled with Runge-Kutta Fehlberg integration scheme is employed to tackle the problem in section two. Both numerical and graphical results are presented and

discussed quantitatively with respect to various parameters embedded in the system in section three.

## 2. Model Formulation

Consider the hydromagnetic turbulent flow of an incompressible electrically conducting fluid past a moving vertical plate in a rotating system with thermal radiation in the presence of a uniform transverse magnetic field of strength  $B_0$ . Choose the coordinate system in such a way that  $x$ -axis is along the plate in upward direction,  $z$  - axis normal to plane of the plate in the fluid and  $y$  - axis perpendicular to  $xz$ - plane. The uniform transverse magnetic field  $B_0$  is applied in a direction which is parallel to  $z$  - axis. It is assumed that no applied or polarization voltages exist so electric field is zero. The magnetic Reynolds number is very small and the induced magnetic field produced by fluid motion is negligible in comparison to applied one. The plate moves along  $x$ -axis with velocity  $U_0$  and both the plate and the fluid rotate in unison with uniform angular velocity  $\Omega$  about  $z$  - axis as shown in figure 1.

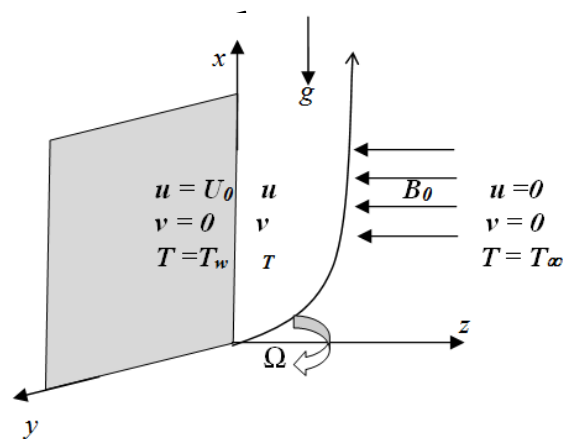


Figure 1. Geometry of the problem

Taking into consideration the assumptions made above together with the application of Reynolds averaged Navier-Stokes (or RANS) methods, the governing equations for hydromagnetic turbulent boundary layer flow of a viscous, incompressible, electrically conducting fluid with Hall effects and thermal radiations, under Boussinesq approximation, in a rotating frame of reference are given by [10, 14, 18, 20]

$$-2\Omega\bar{v} = \nu \frac{d^2\bar{u}}{dz^2} + \frac{d}{dz} \left( n^2 z^2 \left( \frac{d\bar{u}}{dz} \right)^2 \right) + g\beta(\bar{T} - T_\infty) - \frac{\sigma B_0^2}{\rho} \left( \frac{\bar{u} - m\bar{v}}{1 + m^2} \right) \quad (1)$$

$$2\Omega\bar{u} = \nu \frac{d^2\bar{v}}{dz^2} + \frac{d}{dz} \left( n^2 z^2 \left( \frac{d\bar{v}}{dz} \right)^2 \right) - \frac{\sigma B_0^2}{\rho} \left( \frac{\bar{v} + m\bar{u}}{1 + m^2} \right) \quad (2)$$

$$\alpha \frac{d^2 \bar{T}}{dz^2} + \frac{d}{dz} \left( \frac{n^2 z^2}{Pr_t} \frac{d\bar{u}}{dz} \frac{d\bar{T}}{dz} \right) - \frac{1}{\rho c_p} \frac{dq_r}{dz} + \frac{v}{c_p} \left[ \left( \frac{d\bar{u}}{dz} \right)^2 + \left( \frac{d\bar{v}}{dz} \right)^2 \right] + \frac{\sigma B_0^2}{\rho c_p} \left[ \frac{(\bar{u} - m\bar{v})^2 + (\bar{v} + m\bar{u})^2}{(1+m^2)^2} \right] = 0, \quad (3)$$

with

$$\bar{u} = U_0, \quad \bar{v} = 0, \quad \bar{T} = T_w, \quad \text{at } z = 0, \quad (4)$$

$$\bar{u} = 0, \quad \bar{v} = 0, \quad \bar{T} = T_\infty, \quad \text{as } z \rightarrow \infty, \quad (5)$$

heat at constant pressure,  $\beta$  is the thermal expansion coefficient and  $Pr_t$  is the turbulence Prandtl number. It is assumed that the fluid medium is optically thin with relatively low density and following Cogley *et al.* [13], the radiative heat flux  $q_r$  is given as

$$\frac{\partial q_r}{\partial z} = -4\varepsilon^2 (\bar{T} - T_\infty), \quad (6)$$

where  $\varepsilon \ll 1$  is the radiation absorption coefficient. Moreover, it is important to note that in order to obtain the turbulent flow model equations (1)-(3), we have decomposed the flow variables  $u$ ,  $v$  and  $T$  into time-mean and fluctuating part, i.e.

$$u = \bar{u} + u', \quad v = \bar{v} + v', \quad T = \bar{T} + T', \quad \bar{u} = \frac{1}{t} \int_0^t u dt, \quad (7)$$

$$\bar{v} = \frac{1}{t} \int_0^t v dt, \quad \bar{T} = \frac{1}{t} \int_0^t T dt,$$

and apply the time-averaging on time dependent momentum and energy balance equations under laminar flow scenario such that the time averaging of fluctuating part (with prime symbol) is zero. Both Boussinesq approximation and Prandtl mixing length hypothesis [20] are also applied to resolve the time-averaging of product of fluctuating part i.e.

$$\frac{\partial \bar{u}}{\partial t} = \frac{\partial \bar{v}}{\partial t} = \frac{\partial \bar{T}}{\partial t} = 0, \quad \overline{u'w'} = -n^2 z^2 \left( \frac{\partial \bar{u}}{\partial z} \right)^2, \quad \overline{v'w'} = -n^2 z^2 \left( \frac{\partial \bar{v}}{\partial z} \right)^2, \quad (8)$$

$$\overline{T'w'} = -\varepsilon_H \frac{\partial \bar{T}}{\partial z}, \quad \varepsilon_H = \frac{\varepsilon_m}{Pr_t},$$

$$\varepsilon_m = n^2 z^2 \frac{\partial \bar{u}}{\partial z}, \quad \frac{\partial^2}{\partial x^2} \ll \frac{\partial^2}{\partial z^2}, \quad \frac{\partial \bar{u}'^2}{\partial x} \ll \frac{\partial \bar{u}'w'}{\partial z}, \quad \frac{\partial \bar{v}'^2}{\partial x} \ll \frac{\partial \bar{v}'w'}{\partial z},$$

$$\frac{\partial \bar{u}'T'}{\partial x} \ll \frac{\partial \bar{T}'w'}{\partial z},$$

where the terms  $\varepsilon_M$  and  $\varepsilon_H$  are known as the momentum eddy diffusivity and the thermal eddy diffusivity. We introduce the following dimensionless variables and parameters into equations (1)-(5);

$$\eta = \frac{zU_0}{\nu}, \quad U = \frac{\bar{u}}{U_0}, \quad V = \frac{\bar{v}}{U_0}, \quad \theta = \frac{\bar{T} - T_\infty}{T_w - T_\infty}, \quad M = \frac{\sigma B_0^2 \nu}{\rho U_0^2}, \quad m = w_e \tau_e, \quad (9)$$

$$R = \frac{\Omega \nu}{U_0^2}, \quad Pr = \frac{\nu}{\alpha}, \quad Ec = \frac{U_0^2}{c_p (T_w - T_\infty)}, \quad Gr = \frac{g \beta \nu (T_w - T_\infty)}{U_0^3},$$

$$Nr = \frac{4\varepsilon^2 \nu}{\rho c_p U_0^2},$$

and obtain

$$-2RV = \frac{d^2 U}{d\eta^2} + \frac{d}{d\eta} \left( n^2 \eta^2 \left( \frac{dU}{d\eta} \right)^2 \right) + Gr\theta - M \left( \frac{U - mV}{1+m^2} \right) \quad (10)$$

$$2RU = \frac{d^2 V}{d\eta^2} + \frac{d}{d\eta} \left( n^2 \eta^2 \left( \frac{dV}{d\eta} \right)^2 \right) - M \left( \frac{V + mU}{1+m^2} \right) \quad (11)$$

$$\frac{1}{Pr} \frac{d^2 \theta}{d\eta^2} + \frac{d}{d\eta} \left( \frac{n^2 \eta^2}{Pr_t} \frac{dU}{d\eta} \frac{d\theta}{d\eta} \right) + Nr\theta + Ec \left[ \left( \frac{dU}{d\eta} \right)^2 + \left( \frac{dV}{d\eta} \right)^2 \right] + MEc \left[ \frac{(U - mV)^2 + (V + mU)^2}{(1+m^2)^2} \right] = 0 \quad (12)$$

with

$$U = 1, \quad V = 0, \quad \theta = 1, \quad \text{at } \eta = 0, \quad (13)$$

$$U = 0, \quad V = 0, \quad \theta = 0, \quad \text{as } \eta \rightarrow \infty, \quad (14)$$

where  $M$  is the magnetic field parameter,  $Nr$  is the radiation parameter,  $R$  is the rotational parameter,  $Ec$  is the Eckert number,  $Gr$  is the Grashof number and  $Pr$  is the Prandtl number. Following Marchello and Toor [19], for high turbulence intensity, the turbulent Prandtl number is given in term of the fluid Prandtl number as  $Pr_t = \sqrt{Pr}$ . Other quantities of interest are the skin friction ( $C_f$ ) coefficient and the Nusselt number ( $Nu$ ) at the plate surface which are given as

$$C_{fx} = \frac{\tau_{wx}}{\rho U_0^2} = -\frac{dU}{d\eta} \Big|_{\eta=0}, \quad C_{fy} = \frac{\tau_{wy}}{\rho U_0^2} = -\frac{dV}{d\eta} \Big|_{\eta=0},$$

$$Nu = \frac{q_w}{k(T_w - T_\infty)} = -\frac{d\theta}{d\eta} \Big|_{\eta=0}, \quad (15)$$

$$\text{where } \tau_{wx} = -\mu \frac{du}{dz} \Big|_{z=0}, \quad \tau_{wy} = -\mu \frac{dv}{dz} \Big|_{z=0} \quad \text{and} \quad q_w = -k \frac{dT}{dz} \Big|_{z=0}.$$

The nonlinear dimensionless model equations (10)-(13) together with the boundary conditions in (14)-(15) constitute a boundary value problem (BVP) and are solved numerically using shooting method coupled with Runge-Kutta Fehlberg integration scheme. This method involves transforming equations (10)-(15) into a set of initial value problems (IVP) which contain unknown initial values that need to be determined by guessing, after which Runge-Kutta Fehlberg integration scheme is employed to integrate the set of initial value problems until the given boundary conditions are satisfied. The entire computation procedure

is implemented on computer using a program written in MAPLE language. From the process of numerical computation, the mean velocity and mean temperature are obtained. The numerical values are used to compute the skin friction coefficient and the Nusselt number as given in equation (15).

### 3. Results and Discussion

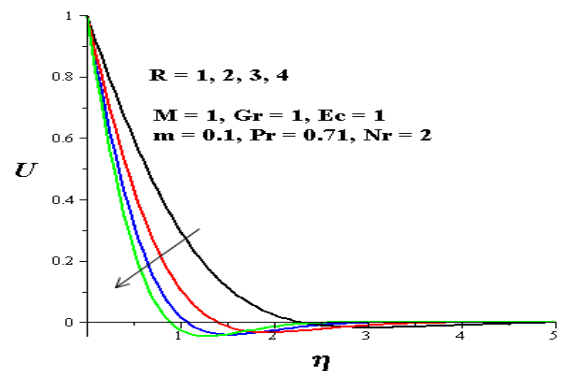
In order to understand the physical situation of the problem and effects of various parameters controlling the flow regime, we have computed the numerical values of mean velocities, mean temperature, the skin friction and the Nusselt number with respect to each parameter variation as shown in figures 2-31. It is important to note that  $Gr > 0$  corresponds to cooling of the plate by convection current while  $Gr < 0$  implies heating of the plate by convection current. Moreover, the cooling problem is often encountered in engineering applications such as cooling of electronic components and nuclear reactors. Numerical values in table 1 depict the effect of parameter variation on skin friction and Nusselt number. From the table, it is observed that the primary skin friction at the plate surface increases with an increase in parameter values of  $M, R, Nr, Pr, Gr < 0$  and decreases with an increase in parameter  $m, Ec$  and  $Gr > 0$ . The secondary skin friction increases with parameter  $m, R, Ec, Gr > 0$  and decreases with parameter  $M, Nr, Pr, Gr < 0$ . An increase in parameter  $m, Nr$  and  $Pr$  increases the Nusselt number while an increase in parameter values of  $M, R, Ec$  and  $Gr$  decreases the Nusselt number.

**Table 1.** Computation showing the skin friction and Nusselt number

$M$	$m$	$R$	$Gr$	$Ec$	$Nr$	$Pr$	$C_{fx}$	$C_{fy}$	$Nu$
1.0	0.1	1.0	0.1	0.1	1.0	0.71	1.21183514	0.81877576	0.83925868
2.0	0.1	1.0	0.1	0.1	1.0	0.71	1.51494067	0.70318128	0.82053306
3.0	0.1	1.0	0.1	0.1	1.0	0.71	1.78458544	0.63077322	0.80364695
1.0	0.5	1.0	0.1	0.1	1.0	0.71	1.21244422	0.94064112	0.84088432
1.0	1.0	1.0	0.1	0.1	1.0	0.71	1.14983533	1.02958701	0.84583895
1.0	0.1	2.0	0.1	0.1	1.0	0.71	1.55777192	1.28916619	0.82453955
1.0	0.1	3.0	0.1	0.1	1.0	0.71	1.84502594	1.62487996	0.80871950
1.0	0.1	1.0	0.5	0.1	1.0	0.71	1.04984232	0.86729541	0.83701858
1.0	0.1	1.0	1.0	0.1	1.0	0.71	0.84543532	0.92903706	0.83237814
1.0	0.1	1.0	0.5	0.1	1.0	0.71	1.45275334	0.74715631	0.84034948
1.0	0.1	1.0	1.0	0.1	1.0	0.71	1.65199584	0.68831179	0.83930426
1.0	0.1	1.0	0.1	0.5	1.0	0.71	1.20895408	0.82012949	0.52261924
1.0	0.1	1.0	0.1	1.0	1.0	0.71	1.20534494	0.82182793	0.12720102
1.0	0.1	1.0	0.1	0.1	2.0	0.71	1.21843079	0.81542480	1.41337547
1.0	0.1	1.0	0.1	0.1	3.0	0.71	1.22393362	0.81296122	2.06881414
1.0	0.1	1.0	0.1	0.1	1.0	4.00	1.23139532	0.81022002	3.49966374
1.0	0.1	1.0	0.1	0.1	1.0	7.10	1.23772125	0.80862109	6.20529525

#### 3.1. Effects of Parameter Variation on Mean Velocities

Figures 2-8 illustrate the effects of parameter variation on the primary mean velocity component. Generally the primary mean velocity is highest at the moving plate surface and tends zero free stream value far away from the plate surface satisfying the prescribed boundary conditions. An increase in the fluid rotation, magnetic field intensity, hall current, thermal radiation and Prandtl number decreases the momentum boundary layer thickness as shown in figures 2-6, consequently the primary mean velocity decreases. In figure 7, it is observed that the momentum boundary layer thickness increases with cooling of the plate by convective current ( $Gr > 0$ ) due to buoyancy force and decreases with heating of the plate by convective current ( $Gr < 0$ ). As the viscous heating increases ( $Ec$ ), the primary mean velocity increases leading to a rise in momentum boundary layer thickness as shown in figure 8. Figures 9-15 depict the secondary mean velocity for some selected values of the parameters. Interestingly, flow reversal is generally observed with its peak value within the boundary layer regime and zero secondary mean velocity at both the plate surface and free stream satisfying the prescribed boundary conditions. As the values of parameter  $R, Nr, M$  and  $Pr$  increase due to combined effects of fluid rotation, magnetic field and thermal radiation, the fluid flows toward the plate surface with a decrease in the reverse flow intensity as shown in figures 9-13. This is expected, since increase in magnetic field enhances the effect of Lorentz force which acts as a resistance to the flow, leading to a decrease in the momentum boundary layer thickness. Figure 14 shows that the secondary mean flow reversal and its peak values increase with cooling ( $Gr > 0$ ) and decreases with heating ( $Gr < 0$ ) of the plate surface by convective current due to buoyancy force. Increase viscous heating ( $Ec$ ) increases the secondary mean flow reversal and the momentum boundary layer thickness as shown figure 15.



**Figure 2.** Primary velocity profiles with increasing  $R$ .

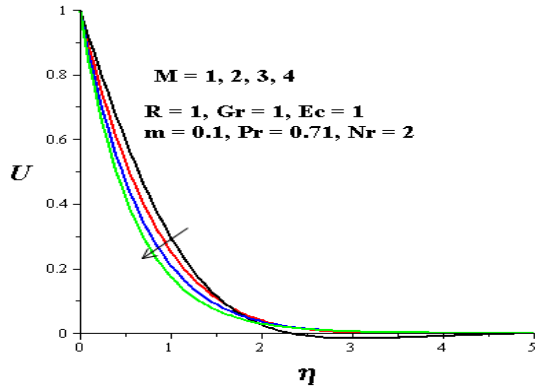


Figure 3. Primary velocity profiles with increasing  $M$ .

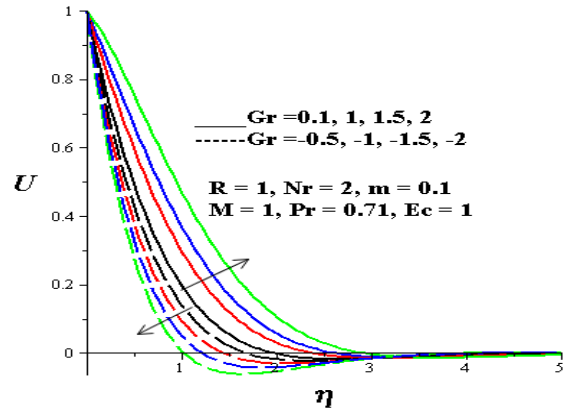


Figure 7. Primary velocity profiles with increasing  $Gr$ .

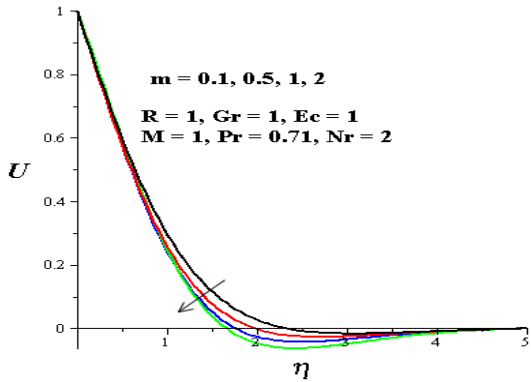


Figure 4. Primary velocity profiles with increasing  $m$ .

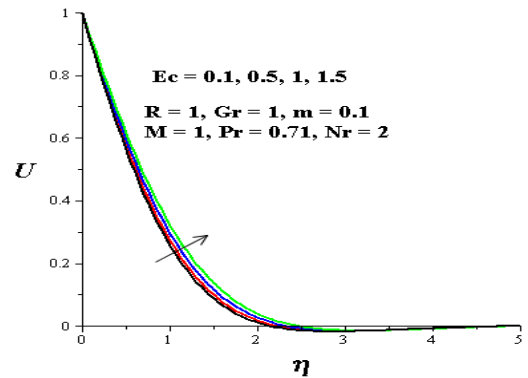


Figure 8. Primary velocity profiles with increasing  $Ec$ .

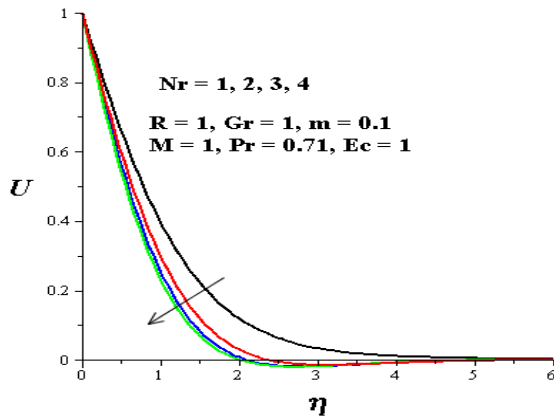


Figure 5. Primary velocity profiles with increasing  $Nr$ .

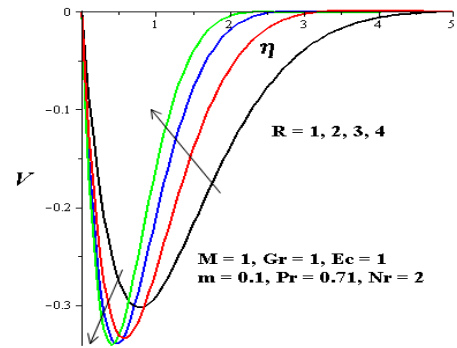


Figure 9. Secondary velocity profiles with increasing  $R$ .

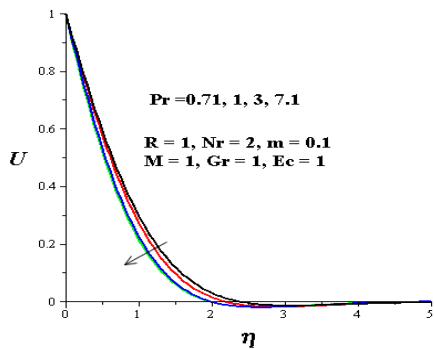


Figure 6. Primary velocity profiles with increasing  $Pr$ .

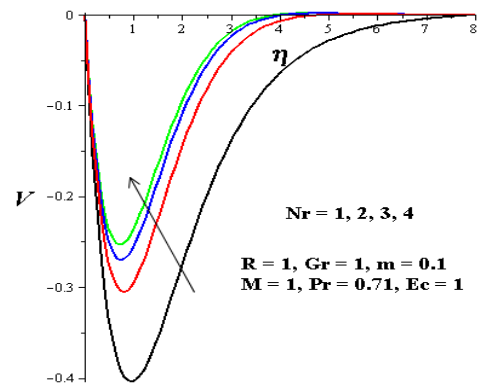


Figure 10. Secondary velocity profiles with increasing  $Nr$ .

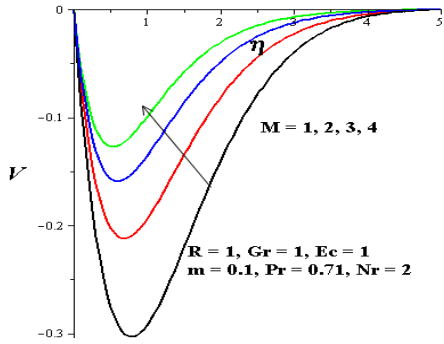


Figure 11. Secondary velocity profiles with increasing  $M$ .

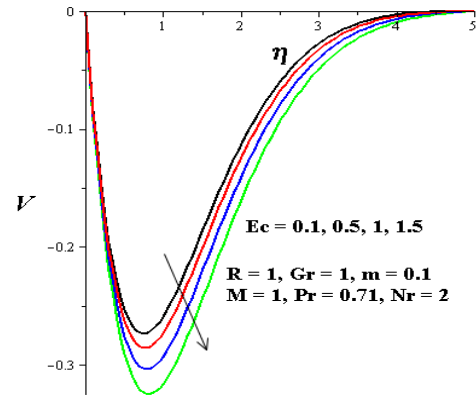


Figure 15. Secondary velocity profiles with increasing  $Ec$ .

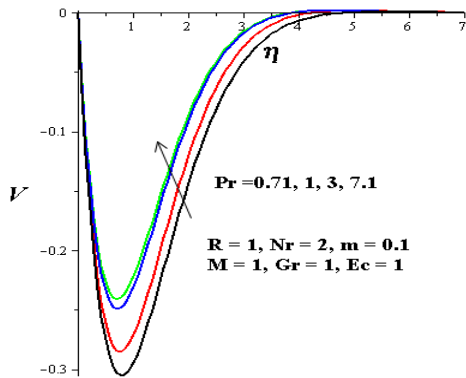


Figure 12. Secondary velocity profiles with increasing  $Pr$ .

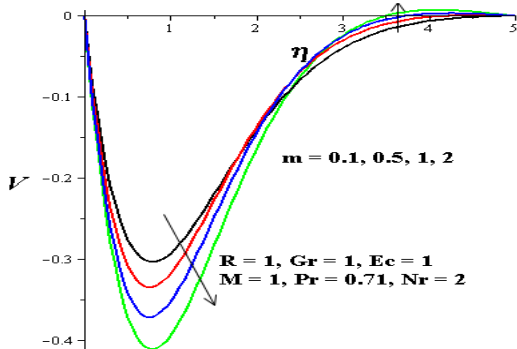


Figure 13. Secondary velocity profiles with increasing  $m$ .

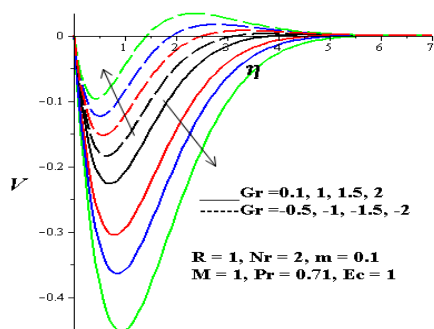


Figure 14. Secondary velocity profiles with increasing  $Gr$ .

### 3.2. Effects of Parameter Variation on Mean Temperature Profiles

Figures 16-22 illustrate the effects of various thermophysical parameters on the mean temperature profiles. It is noteworthy that the mean temperature decreases from the plate surface to the prescribed free stream zero value far away the plate. Meanwhile, as the parameter values of  $R$ ,  $M$  and  $Ec$  increase, the thermal boundary layer thickness increases near the plate surface as shown in figures 16-18. This may be attributed to the combined effects of fluid rotation, viscous and Ohmic heating, and leading to a rise in the mean temperature. In figures 19-21, a decrease in the thermal boundary layer thickness is observed with increasing parameter values of  $Nr$ ,  $m$  and  $Pr$ , consequently, the mean temperature decreases. Figure 22 shows that the mean temperature within the boundary layer regime is enhanced by plate cooling while the heating of the plate by buoyancy force decreases the boundary layer thickness. This can be explained by the fact that the heat is transferred from the plate to the fluid by buoyancy force during cooling leading to a rise in the mean temperature

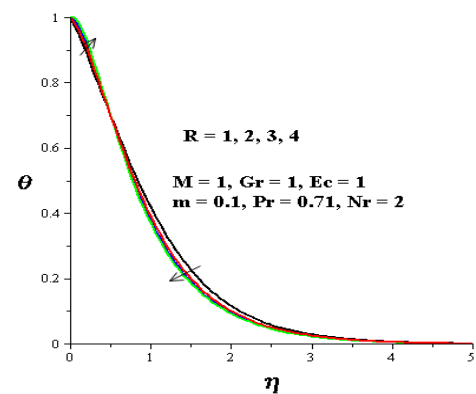


Figure 16. Temperature profiles with increasing  $R$ .

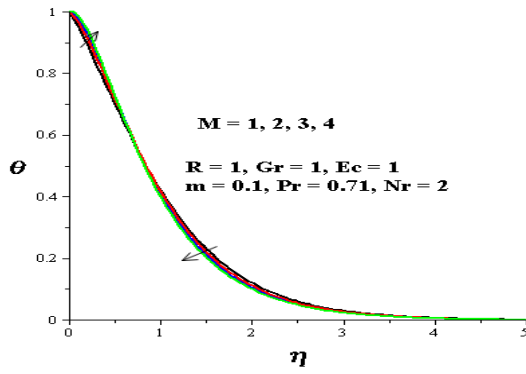


Figure 17. Temperature profiles with increasing  $M$ .

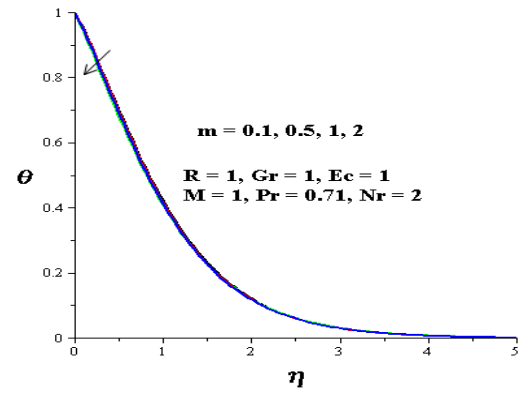


Figure 20. Temperature profiles with increasing  $m$ .

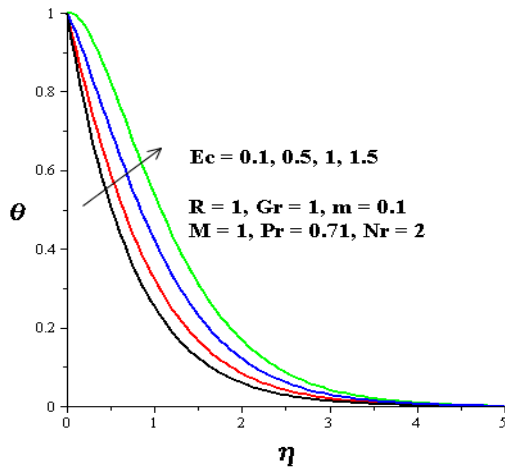


Figure 18. Temperature profiles with increasing  $Ec$ .

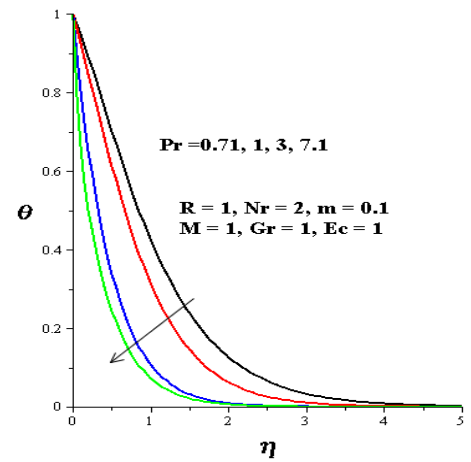


Figure 21. Temperature profiles with increasing  $Pr$ .

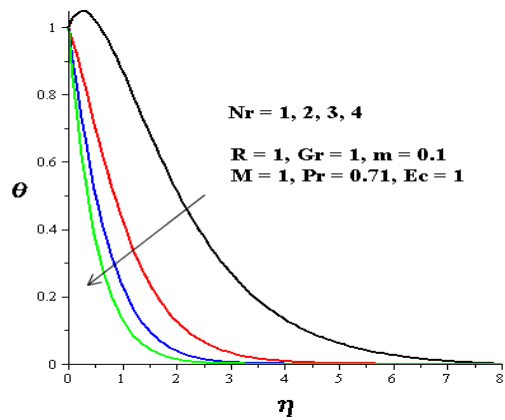


Figure 19. Temperature profiles with increasing  $Nr$ .

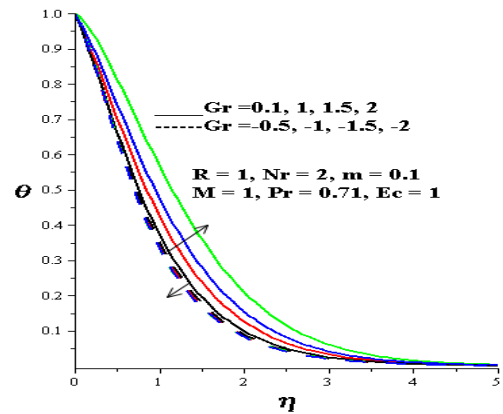


Figure 22. Temperature profiles with increasing  $Gr$ .

### 3.3. Effects of Parameter Variation on Skin Friction and Nusselt Number

Figures 23-25 demonstrate the effects of increasing each thermophysical parameter on the primary skin friction. We observed that the primary skin friction increases with an increase in magnetic field intensity, fluid rotation, thermal radiation, Prandtl number and plate heating by convective current. This can be attributed to a rise in mean velocity gradient at the plate surface. Meanwhile a fall in the primary skin friction is noticed with increasing intensity of

Hall current, viscous heating and plate cooling due to buoyancy force. Figures 26-28 depict the effects of parameter variation on secondary skin friction. A combined increase in the fluid rotation, Hall current, viscous heating and plate cooling due to buoyancy force increases the secondary skin friction while an increase in the Ohmic heating, thermal radiation and Prandtl number decreases the secondary skin friction at the plate surface. In figures 29-31, the variation in Nusselt number is illustrated with different parameters. An increase in Hall current, thermal radiation and Prandtl number enhances the Nusselt number due to a rise in the temperature gradient at the plate surface. However, the Nusselt number decreases with a combined increase in Ohmic and viscous heating, fluid rotation and buoyancy force.

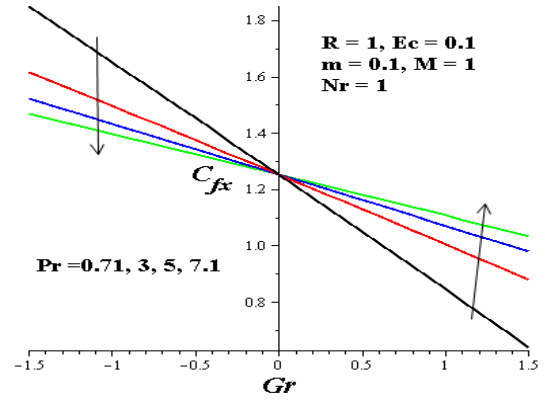


Figure 25. Primary skin friction with increasing Pr.

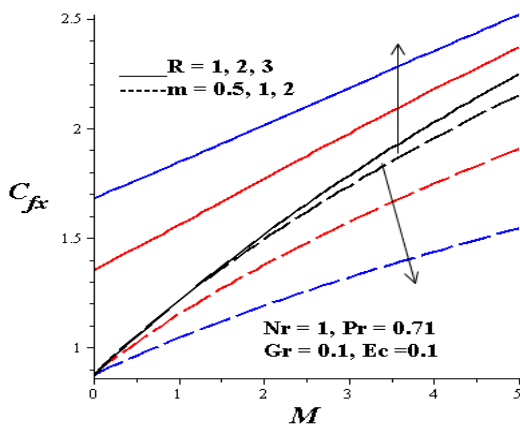


Figure 23. Primary skin friction with increasing M, R and m.

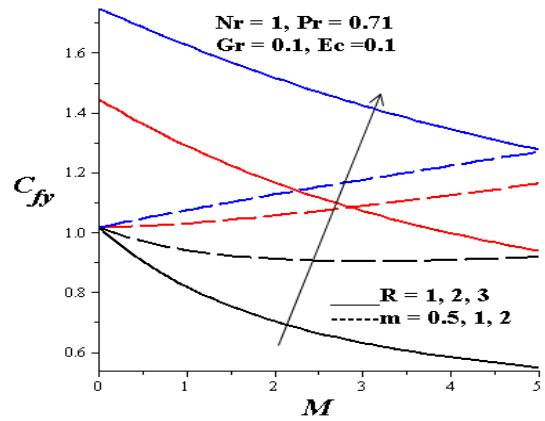


Figure 26. Secondary skin friction with increasing M, R and m.

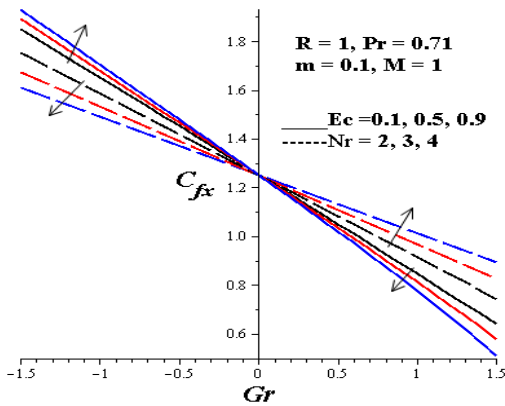


Figure 24. Primary skin friction with increasing Ec, Nr and Gr.

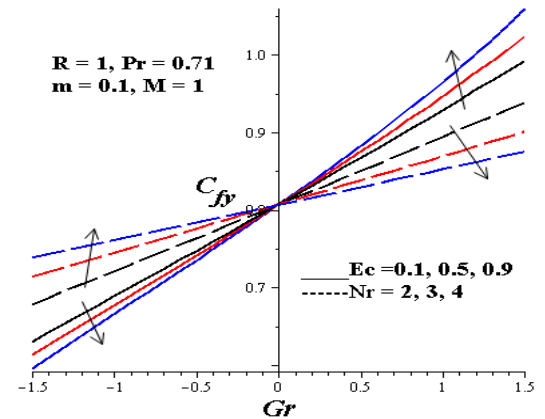


Figure 27. Secondary skin friction with increasing Ec, Nr and Gr.



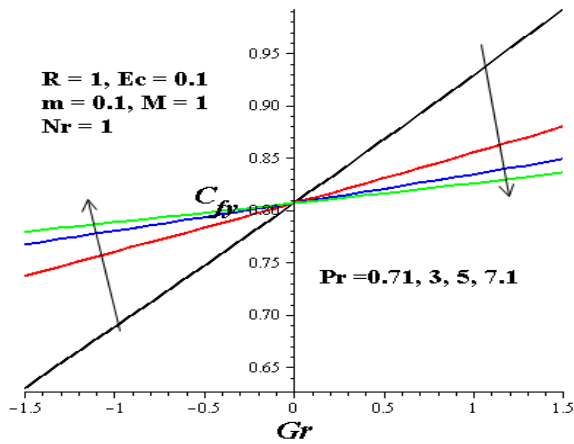


Figure 28. Secondary skin friction with increasing Pr.

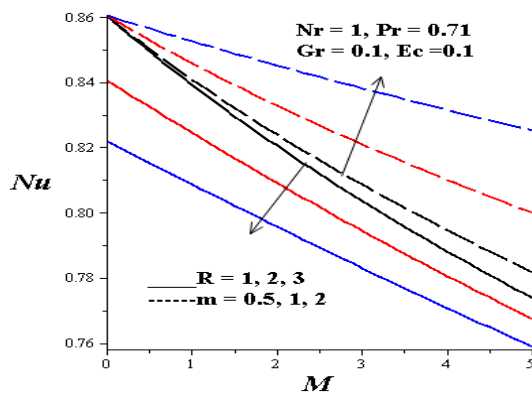


Figure 29. Nusselt number with increasing M, R and m.

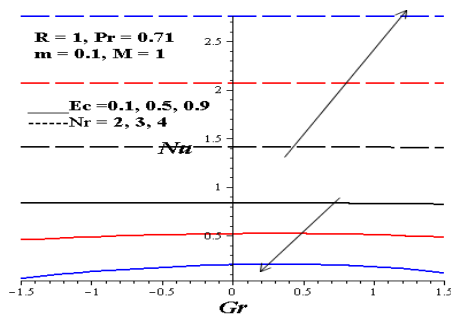


Figure 30. Nusselt number with increasing Ec, Nr and Gr.

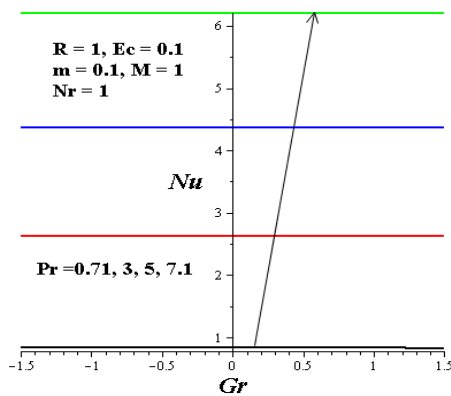


Figure 31. Nusselt number with increasing Pr.

### 4. Conclusions

The hydromagnetic turbulent flow of a conducting fluid over a moving plate in a rotating system with thermal radiation, viscous and Ohmic heating is numerically investigated. The effects of various thermophysical parameters on the mean velocity, mean temperature, skin friction and Nusselt number were obtained. Our results can be summarized as follows:

- The primary mean velocity and momentum boundary layer increases with  $Gr > 0$ ,  $Ec$  and decreases with  $R, M, m, Nr, Pr, Gr < 0$ .
- The secondary mean velocity shows flow reversal and momentum boundary layer increases with  $m, Gr > 0, Ec$  and decreases with  $R, Nr, M, Pr, Gr < 0$ .
- The mean temperature and thermal boundary layer increases with  $R, M, Ec, Gr > 0$  and decreases with  $Nr, Pr, m, Gr < 0$ .
- The primary skin friction increases with  $M, R, Nr, Pr, Gr < 0$  and decreases with  $m, Ec$  and  $Gr > 0$ .
- The secondary skin friction increases  $m, R, Ec, Gr > 0$  and decreases  $M, Nr, Pr, Gr < 0$ .
- The Nusselt number increase with  $m, Nr, Pr$  increases and decreases with  $M, R, Ec, Gr$ .

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