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# A Mathematical Aspect of FRW - Cosmological Models Varying $t$ and $\Lambda$

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**Abstract:** In this paper discussing about the mathematical aspect of variables such as cosmological constant ( $\Lambda$ ) and gravitational constant ( $G$ ) under some suitable assumptions,  $G(t) = F \exp(3H\gamma t)$  (where  $F$  is constant and  $\gamma$  is equation of state parameter), and the cosmological constant.  $\Lambda \propto t^{-2}$ . The exact solutions for the field equations and discuss some physical properties of the cosmological model. An expanding universe is found by using a relation between Friedmann equation and an equation of state. The cosmological constant has been involved on several occasions to correct some seemingly real difficulty with the cosmological predictions of standard general relativity. The most notable of these where its initial use by Einstein to produce static universe. The physical situation regarding  $\Lambda$  is unclear. The other cosmological parameters Hubble constant ( $H$ ), the density parameter ( $\rho$ ), Gravitational constant ( $G$ ) etc, may not constant at all. In this work emphasize to investigate these constants vary with time.

**Keywords:** Friedmann Model, Cosmological Constant, Inflation, Hubble Parameter

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## 1. Introduction

“Einstein Universe” is one of Friedmann solutions of Einstein’s field equations for the value of cosmological constant ( $\Lambda$ ). This is only stationary solution of all Friedmann’s solutions, and because it is stationary, it is thought to be non-physical by majority of astronomers. Those astronomers thinks that universe is expanding because there is observed phenomenon of Hubble red shift and it is interpreted by those astronomers as a Doppler’s shift caused by galaxies moving away from our own galaxy. Therefore, it is thought that the real solutions of Einstein’s field equations cannot be stationary. As discussed earlier by many researchers [4, 5] that a constant ( $\Lambda$ ) cannot explain the huge difference between the cosmological constant inferred from observations and energy density resulting from quantum field theories.

The cosmological constant ( $\Lambda$ ) was introduced such that it was large in the early universe and then decayed with evolution. A number of models with different decay laws for the variation of cosmological term were investigated during last two decades by Abdel-Rahaman 1990 [6] and Berman 1991 [2], Maharaj and Naido 1993 [7] who obtained cosmological model with time dependent  $G$  and  $\Lambda$  terms,

and suggested that  $\Lambda \sim t^{-2}$ .

Many authors [1, 3, 5, 6] have proposed linking of the variation of  $G$  with that of  $\Lambda$  within the rules of general relativity. Mathematically well-posed gravitation theories were developed in which Einstein’s theory of general relativity was generalized to include a varying  $\Lambda$  by deriving it from Friedmann equation satisfying a conservation equation. There has been considerable interest in solutions of Einstein’s equations of early universe both from cosmological as well as mathematical point of view [1, 5, 6, 7].

A viscous model with variable gravitational and cosmological constant has been considered. The role of viscosity in cosmology has been studied by several authors [1-4]. The bulk viscosity associated with grand unified theory phase transition can lead to the inflationary universe scenario. It is well known that in an early stage of the universe when neutrino decoupling occurred, the matter behaves like viscous fluid [8]. The coefficient of viscosity is known to decrease as the universe expands. Beesham [10] studied a universe consisting of a cosmological constant ( $\Lambda \sim t^{-2}$ ) and bulk viscosity. More recently Abdel Rahman considered a model in which the gravitational constant  $G$  varies with time but energy is conserved [5, 6, 9]. In the present work will investigate the effects of viscosity in a universe vary with

time in different cosmological constant like  $H, G, \Lambda$  etc.

## 2. The Model

In a Robertson-Walker universe

$$d\tau^2 = dt^2 - R^2(t) \left[ \frac{dr^2}{1-kr^2} + r^2(d\theta^2 + \sin^2\theta d\varphi^2) \right] \quad (2.1)$$

where  $k$  is the curvature index

Einstein's field equations with time dependent cosmological and gravitational "constants"

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi GT_{\mu\nu} + \Lambda g_{\mu\nu} \quad (2.2)$$

and the perfect fluid energy momentum tensor

$$T_{\mu\nu} = (p + \rho)U_\mu U_\nu - p g_{\mu\nu} \quad (2.3)$$

yield the two independent equations

$$3\frac{\dot{R}}{R} = -4\pi G \left( 3p + \rho - \frac{\Lambda}{4\pi G} \right) \quad (2.4)$$

$$3\frac{\dot{R}^2}{R^2} = 8\pi G \left( \rho + \frac{\Lambda}{8\pi G} \right) - \frac{3k}{R^2} \quad (2.5)$$

Without the cosmological constant solution of equation (2.4) takes the space-space component of Friedmann equation as

$$\frac{\ddot{R}}{R} = -\frac{4\pi G}{3}(3p + \rho) \quad (2.6)$$

and equation(2.5) takes the time-time component of the form

$$\left( \frac{\dot{R}}{R} \right)^2 = \frac{8\pi G\rho}{3} - \frac{k}{R^2} \quad (2.7)$$

For flat universe  $k=0$  and  $\rho=0$ , when the matter is not considered as ideal gas in the early universe. Then the equation (2.5) takes the form

$$H = \frac{\dot{R}}{R} = \sqrt{\frac{\Lambda}{3}}$$

where  $H=\sqrt{\Lambda/3}$  is the Hubble Parameter.

Using this statement we get de-sitter universe

$$R(t) = R_0 e^{Ht} \quad (2.8)$$

This is also the inflationary solution.

From equation (2.4) and (2.5) gives

$$3(p + \rho)\dot{R} = -\left( \frac{\dot{G}}{G}\rho + \dot{\rho} + \frac{\dot{\Lambda}}{8\pi G} \right)R \quad (2.9)$$

The conservation of energy and momentum yields

$$3(p + \rho) = -R \frac{d\rho}{dR} \quad (2.10)$$

The effect of the bulk viscosity in the field equation is to

replace  $p$  by  $p - \eta H$ , where  $\eta$  is the viscosity co-efficient.

$$3(p - \eta H + \rho)\dot{R} = -\left( \frac{\dot{G}}{G}\rho + \dot{\rho} + \frac{\dot{\Lambda}}{8\pi G} \right)R$$

$$3(p + \rho)\dot{R} - 3\eta H\dot{R} = -\left( \frac{\dot{G}}{G}\rho + \frac{\dot{\Lambda}}{8\pi G} \right)R - \dot{\rho}R$$

It follows immediately that

$$3\eta H\dot{R} = \left( \frac{\dot{G}}{G}\rho + \frac{\dot{\Lambda}}{8\pi G} \right)R \quad (2.11)$$

and

$$\dot{\rho} + 3H(p + \rho) = 0 \quad (2.12)$$

Equation (2.11) can also be written as

$$\frac{3\eta H}{R} = \rho \frac{G'}{G} + \frac{\Lambda'}{8\pi G} \quad (2.13)$$

Where prime denotes derivative with respect to scale factor  $R$ , while dot is the derivative with respect to  $t$ . For considering a flat universe,  $k=0$

Equation (2.5) leads to

$$8\pi G\rho = 3H^2(1 - \beta) \quad (2.14)$$

Where

$$\beta = \frac{\Lambda}{3H^2} \quad (2.14a)$$

and  $\beta = \text{Constant}$

Also the equation of state is

$$p = (\gamma - 1)\rho \quad (2.15)$$

Where  $1 \leq \gamma \leq 2$

In equation (2.12) and (2.15) leads to

$$\rho = AR^{-3\gamma} \quad (2.16)$$

where  $A$  is the constant of integration.

This is the general solution of density of the universe. If  $\gamma=1$ , then  $\rho \propto 1/R^3$  is the dust like universe or matter dominant universe [11] and if  $\gamma=4/3$ , then  $\rho \propto 1/R^4$  is the radiation dominant universe. If  $\gamma=2$ , then  $\rho \propto 1/R^6$  for stiff fluid.

$$\rho_d \propto R^{-3}, \rho_r \propto R^{-4}, \rho_d/\rho_r \propto R$$

To distinguish dust and radiation, we will use the symbols  $\rho_d$  and  $\rho_r$  to denote their respective densities. Thus in an expanding universe the dust begins to dominate over radiation as the expansion continuous beyond a certain limit. The evidence for 'dust' in the form of visible matter in the galaxies, intergalactic medium etc., suggests that at present

$$\rho_d \cong 3 \times 10^{-31} \text{ gcm}^{-3} \quad (2.17)$$

From the observational evidence which suggests that the universe may very likely contain dark matter far in excess of the visible matter. It is possible in view of this that the actual matter density is higher. Let us therefore enhance by a dark matter factor,  $\alpha > 1$  and write

$$\rho_d = 3 \times 10^{-31} \text{ gm cm}^{-3} \quad (2.18)$$

The radiation background exists in all observable wavelengths but the most dominant one is in microwaves in the form of a blackbody radiation of temperature  $\sim 2.7$  K, this corresponds to an energy density

$$\rho_r c^2 \approx 4 \times 10^{-13} \text{ erg cm}^{-3} \quad (2.19)$$

Thus at present

$$\frac{\rho_d}{\rho_c} = \frac{3\alpha}{4} \times 10^3 \text{ i.e. } \rho_d \gg \rho_r$$

As the universe expands this ratio will further increase, not only is the universe matter dominated today, it will continue that way so long as it expands. Clearly for a small enough value of  $R$ ,  $\rho_d$  was less than  $\rho_r$ . From equation  $\frac{\rho_d}{\rho_c} \propto R$  and

$\frac{\rho_d}{\rho_c} \approx \frac{3\alpha}{4} \times 10^3$  we see that this would occur at an epoch of red shift

$$z_{eq} \approx \frac{3\alpha}{4} \times 10^3 \quad (2.20)$$

at  $R_{eq} = \frac{R_0}{1+z_{eq}}$

The matter and radiation densities were comparable. For  $z > z_{eq}$  and  $R < R_{eq}$ , the universe was radiation dominated.

Now from equation (2.14) we have

$$G = \frac{3(1-\beta)H^2}{8\pi\rho} \quad (2.21)$$

Again from equation (2.16) we get

$$\rho = AR_0^{-3\gamma} \exp(-3H\gamma t)$$

Putting the value in equation (2.21) we get finally,

$$G(t) = F \exp(3H\gamma t)$$

where  $F = \frac{3H^2(1-\beta)}{8\pi A} R_0^{3\gamma} = \text{constant}$

This is another solution of the Gravitational constant [12] which is not a constant at all but a function of  $t$ .

From equation (2.5) and  $\Lambda = 0$ ,  $k = 0$  and  $\rho \propto R^{-3}$  and  $\rho \propto R^{-4}$ , for matter and radiation dominant universe respectively,

$$R(t) \propto t^{\frac{2}{3}} \text{ and } R(t) \propto t^{\frac{1}{2}}$$

From equation (2.5) for  $\Lambda = 0$ ,  $k = 0$  and  $\rho = 0$  then

$$H = \frac{\dot{R}}{R} \text{ for } R(t) = kt^{\frac{2}{3}} \text{ and } \dot{R}(t) = k \frac{2}{3} t^{-\frac{1}{3}}$$

And from equation (2.14a) we get

$$\Lambda(t) \propto t^{-2}$$

Which means the cosmological constant is the function of time and also the inverse square of time. From the above, it can be easily seen that energy density ( $\rho$ ) and cosmological constant ( $\Lambda$ ) are decreasing while gravitational constant ( $G$ ) is also decreasing during the expansion of the universe.

### 3. Discussions and Conclusion

This paper to present a variety of classes of solutions to Einstein's field equations containing different parameters  $H, \Lambda, G$ , and bulk viscosity. We have listed the forms of  $R, \rho, G$  and  $\Lambda$  as functions of time. It is clear that these quantities depend on the dimensionality of space-time. The cosmological term  $\Lambda$  decides the behavior of the universe in the model. In this paper we have obtained a negative value of  $\Lambda$  which will correspond to positive effect mass density. Hence we can get a universe that expands and then reconstruct. The observations on magnitude and red shift that our universe may be an accelerating one or otherwise with induced cosmological density, through the cosmological  $\Lambda$  term [11, 12]. These solutions turned out to be power law, as are these suggested by Bermann [2]. It was shown by Maharaj [7] that this is also valid for the scalar tensor theory of law  $k = 0$ . All power law solutions exhibit the feature of viscous solutions with variable with  $G$ . We have found that  $\Lambda \sim t^{-2}$  for all epochs except the inflationary one.

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